# Mathematical Study on MHD Squeeze Flow between Two Parallel Disks with Suction or Injection via HAM and HPM and Its Applications

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Abstract-In this paper, we are considering the problem of magneto-hydrodynamic MHD squeeze flow of an electrically conducting fluid between two infinite, parallel disks are investigated. The analytical method called Homotopy Analysis Method (HAM) and Homotopy Perturbation Method (HPM) are used to compute an approximation for the solution of nonlinear differential equations governing on the problem. The results of the mentioned methods are compared with a type of numerical analysis as Boundary Value Problem method.

**Keywords** — Magneto-hydrodynamic Homotopy Perturbation Method, Squeeze flow, Temprature, nonlinear differential equations, incompressible flow.

#### I. INTRODUCTION

The application of a MHD fluid in lubrication prevents the adverse impact of temperature on the fluid viscosity when the system operates under boundary conditions. The problem considered is of general interest in the theory of magnetohydrodynamic lubrication and other related applications. In particular, the results of the present investigation are directly applicable to the hydrodynamics of high temperature bearings lubricated with liquid metals. A number of theoretical and experimental investigations into magneto-hydrodynamic effects in lubrication have been reported. These include among other works of Hughes and Elco [1], Kuzma et al. [2] and Krieger et al. [3]. Most scientific problems such as twodimensional viscous flow between slowly expanding or contracting walls with weak permeability and other fluid mechanic problems are inherently nonlinear. In most cases, these problems do not admit analytical solution, so these equations should be solved using special techniques. In recent years, much attention has been devoted to the newly developed methods to construct an analytic solution of equation such as the method included the Perturbation techniques. Perturbation techniques are too strongly dependent upon the so-called "small

parameters" [4]. Other different methods have introduced to solve nonlinear equations such as the  $\delta$ -expansion method [5], Adomian's decomposition method [6], Homotopy Perturbation Method (HPM) [7–10], Variational Iteration Method (VIM) [11–14], Homotopy analysis method [15-18], Optimal Homotopy Asymptopic Method (OHAM)[19,20] and optimal Homotopy Perturbation Method (OHPM)[21]. In this letter, analytical solutions of equations arising nonlinear of magnetohydrodynamic MHD squeeze flow of an electrically conducting fluid between two infinite, parallel disks have been studied by the two analytical methods. These methods called Homotopy Analysis Method (HAM), Homotopy Perturbation Method (HPM) do not have small parameters. Obtaining the analytical solution of the models and comparing with the numerical result reveal the capability, effectiveness and convenience of HAM and HPM. These methods give successive approximations of high accuracy solution. Kumar et al. [23]investigated a finite difference technique for reliable MHD steady flow through channels permeable boundaries.Kumaret al [24] investigated MHD free convective fluctuating flow through a porous effect with variable permeability Parameter.Kumar et al. [25] investigated mathematical analysis of MHD on laminar mixed convection of newtonian fluid between vertical parallel plates channel. Kumar et al. [26] investigated a Crank-Nicholson scheme to transient MHD free convective flow through semiinfinite vertical porous plate with constant suction and temperature dependent heat source.

## II. MATHEMATICAL MODEL

In the present investigation, consider an axisymmetric incompressible flow between two parallel infinite disks, which at time  $t^*$ , are space a distance

 $H(1-at^*)^{\frac{1}{2}}$  apart and a magnetic field

proportional to  $B_0(1-at^*)^{-\frac{1}{2}}$  is applied perpendicularly to the disks [6, 22]. The upper disk at  $z = H(1-at^*)^{\frac{1}{2}}$  is moving with velocity  $-\frac{1}{2}\alpha H(1-at^*)^{-\frac{1}{2}}$  towards the stationary lower disk at z = 0. The axial coordinate is denoted by  $z^*$ and the radial coordinate by  $r^*$ . With the axial and

and the radial coordinate by r. With the axial and radial velocities denoted by  $w^*$  and  $u^*$ , respectively, we introduce the following quantities:

$$u^{*} = \frac{\alpha r^{*}}{2(1-at^{*})} f'(\eta), \qquad w^{*} = \frac{\alpha H}{\sqrt{1-at^{*}}} f(\eta), \qquad B = \frac{B_{0}}{\sqrt{1-at^{*}}}$$
$$\eta = \frac{z^{*}}{H\sqrt{1-at^{*}}}, \qquad r = r^{*}, \qquad t = t^{*}$$

The equation of continuity is satisfied and the momentum equations are reduced to:

$$f'''(\eta) - S[\eta f''(\eta) + 3f''(\eta) - 2f(\eta)f'''(\eta)] - M^2 f''(\eta) = 0$$
(2)

Where 
$$\frac{1}{r}\frac{\partial p^*}{\partial r} = p_1(t)$$
 has been used,  $S = \frac{\rho H^2}{2\nu}$ 

and  $M = \frac{\sigma B_0^2}{\mu}$  that  $\rho$  denotes density,  $\upsilon$  denotes

kinetic viscosity and  $\sigma$  denotes fluid electrical conductivity. The boundary conditions are given by:

$$f(0) = A,$$
  $f'(0) = 0,$   $f(1) = \frac{1}{2},$   $f'(1) = 0$  (3)

Where, A is the constant parameter such that A > 0 corresponds to suction and A < 0 to injection.

#### III. APPLICATION OF HOMOTOPY ANALYSIS METHOD

For HAM solutions, we choose the initial guess and auxiliary linear operator in the following form:

$$f_{0}(\eta) = (-1+2A)\eta^{3} + \left(\frac{3}{2} - 3A\right)\eta^{2} + A, \quad (4)$$

$$L(f) = f''', \quad (5)$$

$$L\left(\frac{1}{6}c_{1}y^{3} + \frac{1}{2}c_{2}y^{2} + c_{3}y + c_{4}\right) = 0, \quad (6)$$

Where  $c_i (i = 1, 2, 3, 4)$  are constant. Let  $P \in [0, 1]$  denotes the embedding parameter and  $\hbar$  indicates non –zero auxiliary parameters.

#### Zeroth –order deformation equations

$$(1-P)L[F(\eta,p)-f_0(\eta)] = p\hbar H(\eta)N[F(\eta,p)]$$
(7)

$$F(0; p) = A;$$
  $F'(0; p) = 0,$   $F(1; p) = \frac{1}{2},$   $F'(1; p) = 0$  (8)

$$N[F(\eta;p)] = \frac{d^{4}F(\eta;p)}{d\eta^{4}} + S\left[\eta \frac{d^{3}F(\eta;p)}{d\eta^{3}} + 3\frac{d^{2}F(\eta;p)}{d\eta^{2}}\right] - M^{2} \frac{d^{2}F(\eta;p)}{d\eta^{2}} - 2S\left[F(\eta;p)\frac{d^{3}F(\eta;p)}{d\eta^{3}} - \frac{dF(\eta;p)}{d\eta}\frac{d^{2}F(\eta;p)}{d\eta^{2}}\right]$$
(9)

For 
$$p = 0$$
 and  $p = 1$  we have  
 $F(\eta; 0) = f_0(\eta)$   $F(\eta; 1) = f(\eta)$  (10)

(1)

When p increases from 0 to 1 then  $F(\eta; p)$  varies from  $f_0(\eta)$  to  $f(\eta)$ . By Taylor's theorem and using Eq. (9),  $F(\eta; p)$  can be expanded in a power series of p as follows:

$$F(\eta; p) = f_0(\eta) + \sum_{m=1}^{\infty} f_m(\eta) p^m , \qquad f_m(\eta) = \frac{1}{m!} \frac{\partial^m (F(\eta; p))}{\partial p^m} \bigg|_{p=0}$$
(11)

In which  $\hbar$  is chosen in such a way that this series is convergent at p = 1, therefore we have in according to Eq. (10) that

$$f(\eta) = f_0(\eta) + \sum_{m=1}^{\infty} f_m(\eta),$$
 (12)

$$m^{th} - order \ deformation \ equations$$

$$L[f_m(\eta) - \chi_m \ f_{m-1}(\eta)] = \hbar H(\eta) R_m(\eta) \qquad (13)$$

$$F(0;p) = 0; \quad F'(0;p) = 0, \quad F(1;p) = 0, \quad F'(1;p) = 0 \quad (14)$$

$$R_m(\eta) = f_{m-1}^{m} + S[\eta \ f_{m-1}^{m} + 3f_{m-1}^{m}] - M^2 f_{m-1}^{m} + \sum_{k=0}^{m-1} 2S \ f_{m-1-k} \ f_k^{m} \qquad (15)$$

Now we determine the convergence of the result, the differential equation, and the auxiliary function according to the solution expression. So, let us assume:

$$H(\eta) = 1 \tag{16}$$

We have found the answer by the maple analytic solution device. The first deformation of the solution are presented below :

$$f_{1}(\eta) = -3(-1+2A)\hbar \left(-\frac{1}{210}S + \frac{1}{105}SA\right)\eta^{7} - 3(-1+2A)\hbar \left(-\frac{1}{30}SA + \frac{1}{60}S\right)\eta^{6} \quad (17)$$
  
$$-3(-1+2A)\hbar \left(\frac{1}{60}M^{2} - \frac{1}{15}S\right)\eta^{5} - 3(-1+2A)\hbar \left(-\frac{1}{24}M^{2} + \frac{1}{6}SA + \frac{1}{8}S\right)\eta^{4}$$
  
$$+\hbar \left(-\frac{13}{70}SA - \frac{39}{140}S + \frac{1}{10}M^{2} + \frac{52}{35}SA^{2} - \frac{1}{5}M^{2}A\right)\eta^{3} + \hbar \left(\frac{5}{28}SA + \frac{19}{280}S\right)$$
  
$$-\frac{1}{40}M^{2} - \frac{22}{35}SA^{2} + \frac{1}{20}M^{2}A\right)\eta^{2}$$

The solutions  $f(\eta)$  were too long to be mentioned here, therefore they are shown graphically.

# IV.APPLICATION OF HOMOTOPY PERTURBATION METHOD

In this section, we employ HPM to solve Eq. (2) subject to boundary conditions Eq.(3).We can construct Homotopy function of Eq. (2) as described in [22]:

$$\begin{split} H \left( f \ , \ p \right) &= \left( 1 - P \right) \left[ f''' \ (\eta) - g_0(y) \right] + p \left\{ f'''(\eta) - S \left[ \eta f'''(\eta) + 3 f''(\eta) - 2 f(\eta) f'''(\eta) \right], \\ &- M^2 f''(\eta) \right\} &= 0, \end{split}$$

Where  $p \in [0,1]$  is an embedding parameter. For p = 0 and p = 1 we have:

$$f(\eta, 0) = f_0(\eta) \qquad , \qquad f(\eta, 1) = f(\eta) \tag{19}$$

Note that when p increases from 0 to 1,  $f(\eta, p)$  varies from  $f_0(\eta)$  to  $f(\eta)$ .By substituting:

$$f(\eta) = f_0(\eta) + p f_1(\eta) + p^2 f_2(\eta) + \dots = \sum_{i=0}^n p^i f_i(\eta) , \quad g_0 = 0$$
 (20)

From equation (33) and rearranging the result based on powers of p-terms, we get:  $P^0: f_n^{m'}(n) = 0$ 

$$F': \quad f_0^{-n}(\eta) = 0$$
  

$$f_0(0) = A; \quad f_0'(0) = 0, \quad f_0(1) = \frac{1}{2}, \quad f_0'(1) = 0$$
(21)

- $P^{1}: \quad f_{1}^{'''}(\eta) + S\eta \left( f_{0}^{'''}(\eta) \right) 2Sf_{0}(\eta) \left( f_{0}^{'''}(\eta) \right) + 3S \left( f_{0}^{''}(\eta) \right) M^{2} \left( f_{0}^{''}(\eta) \right) = 0 \quad (22)$  $f_{1}(0) = 0; \quad f_{1}^{'}(0) = 0, \qquad f_{1}(1) = 0, \qquad f_{1}^{'}(1) = 0$
- $P^{2}: \quad f_{2}^{"'}(\eta) 2S f_{1}(\eta) \Big( f_{0}^{"}(\eta) \Big) + 3S \Big( f_{1}^{"}(\eta) \Big) 2Sf_{0}(\eta) \Big( f_{1}^{"'}(\eta) \Big) + S\eta \Big( f_{1}^{"'}(\eta) \Big)$   $-M^{2} \Big( f_{1}^{"}(\eta) \Big) = 0$   $f_{2}(0) = 0; \quad f_{2}'(0) = 0, \quad f_{2}(1) = 0, \quad f_{2}'(1) = 0$ (23)
- $P^{3}: \quad f_{3}^{"''}(\eta) + 3S\left(f_{2}^{"}(\eta)\right) M^{2}\left(f_{2}^{"}(\eta)\right) 2Sf_{0}(\eta)\left(f_{2}^{"'}(\eta)\right) 2Sf_{2}(\eta)\left(f_{0}^{"'}(\eta)\right) \\ 2Sf_{1}(\eta)\left(f_{1}^{"'}(\eta)\right) + S\eta\left(f_{2}^{"'}(\eta)\right) = 0 \\ f_{3}(0) = 0; \qquad f_{3}^{'}(0) = 0, \qquad f_{3}(1) = 0, \qquad f_{3}^{'}(1) = 0$

(24)

Solving Esq. (21) – (24) with boundary conditions, we have (for example S = 0.1, M = 2, A = 1):

$$f_0(\eta) = \eta^3 - 1.50000000\eta^2 + 1,$$
(25)

 $f_1(\eta) = 0.001428571429\eta^7 - 0.0050000\eta^6 + 0.180000\eta^5 - 0.4125000\eta^4$ 

 $+0.2978571429\eta^{3}-0.06178571429\eta^{2}$ 

(26)

 $f_{2}(\eta) = 0.000007792207792\eta^{11} - 0.00004285714286\eta^{10} + 0.0009087301587\eta^{9} - 0.0036607142\eta^{8} + 0.0188153061\eta^{7} - 0.043820238\eta^{6} + 0.03711428571\eta^{5} - 0.004157738100\eta^{4} - 0.005947057248\eta^{3} + 0.0007824906815\eta^{2}$ 

$$\begin{split} f_{3}(\eta) &= 4.99976214310^{-8} \eta^{15} - 3.74982160710^{-7} \eta^{14} + 0.692501942410^{-6} \eta^{13} \\ &- 0.000038502886 \eta^{12} + 0.0002321360544 \eta^{11} - 0.0008178911564 \eta^{10} \\ &+ 0.0019451346370 \eta^{9} - 0.0030426296780 \eta^{8} + 0.002043508165 \eta^{7} \\ &+ 0.0007230726130 \eta^{6} - 0.001236779829 \eta^{5} - 0.00005608490221 \eta^{4} \\ &+ 0.0001191218468 \eta^{3} + 0.0001223151002 \eta^{2} \end{split}$$

The solution of this equation, when  $p \rightarrow 1$ , will be as follows:

$$f(\eta) = \sum_{i=0}^{3} \lim_{p \to 1} p^{i} f_{i}(\eta)$$
(29)  
$$\stackrel{S = 0.01 , M = 5}{= 0.01 , M = 5}$$
$$f'''(0)$$

Fig. 1 The  $\hbar$  - validity for S = 0.01, M = 5 and different value of A.

-1 ħ -0.5



Fig 2. The  $\hbar$  - validity for S = 0.3, A = -1 and different value of M.

As pointed out by Liao, the convergence and rate of approximation for the HAM solution strongly depends on the value of auxiliary parameter  $\hbar$ . The

auxiliary parameter  $\hbar$  provides us with a convenient way to adjust and control the convergency. The range of  $\hbar$  for convergency is obtained according to figs. 1 and 2. For S = 0.01, M = 5 and -7 < A < 7 the ranges  $-0.4 < \hbar < -1.5$ , for S = 0.3, A = -1 and 0 < M < 5 the

ranges  $-0.5 < \hbar < -1.3$ , give suitable value of  $\hbar$  for convergency. Then,  $\hbar = -0.9$  is a suitable value for ranges which is used for the solution.



Fig 3. The comparison between the HAM, HPM and numerical solutions for  $f(\eta)$ ,  $f'(\eta)$ when S = 0.01, M = 5, A = 1.





Fig 4. The comparison between the HAM, HPM and numerical solutions for  $f(\eta), f'(\eta)$  when

S = 0.01, M = 0, A = -1.



Fig 6. The comparison between the HAM, HPM and numerical solutions for  $f(\eta)$ ,  $f'(\eta)$ when S = 1, M = 3, A = 1.

Table 1 The results of HAM, HPM and Numerical methods for $f(\eta)$	) when
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S = 0.4, M = 2, A = 1

		, ,		
HAM	HPM	NUM	Error of HAM	Error of HPM
1.000000000	1.000000000	1.000000000	0.0000000000	0.0000000000
0.996484004	0.996484005	0.996484005	0.0000000005	0.000000006
0.986428373	0.986428377	0.986428373	0.0000000004	0.000000045
0.970559215	0.970559224	0.970559216	0.000000011	0.000000075
0.949588851	0.949588863	0.949588853	0.000000013	0.000000800
0.924219187	0.924219200	0.924219187	0.000000002	0.000000130
0.895145032	0.895145049	0.895145034	0.000000014	0.0000000160
0.863057399	0.863057414	0.863057400	0.0000000002	0.0000000150
0.828646764	0.828646782	0.828646765	0.000000014	0.000000170
	HAM 1.000000000 0.996484004 0.986428373 0.970559215 0.949588851 0.924219187 0.895145032 0.863057399 0.828646764	HAMHPM1.000000001.000000000.9964840040.9964840050.9864283730.9864283770.9705592150.9705592240.9495888510.949588630.9242191870.9242192000.8951450320.8951450490.8630573990.8630574140.8286467640.828646782	HAMHPMNUM1.000000001.000000001.000000000.9964840040.9964840050.9964840050.9864283730.9864283770.9864283730.9705592150.9705592240.9705592160.9495888510.9495888630.9495888530.9242191870.9242192000.9242191870.8951450320.8951450490.8951450340.8630573990.8630574140.8630574000.8286467640.8286467820.828646765	HAMHPMNUMError of HAM1.000000001.000000001.000000000.0000000000.9964840040.9964840050.9964840050.00000000050.9864283730.9864283770.9864283730.0000000040.9705592150.9705592240.9705592160.00000000110.9495888510.9495888630.9495888530.00000000130.9242191870.9242192000.9242191870.0000000020.8951450320.8951450490.8951450340.00000000140.8630573990.8630574140.8630574000.0000000020.8286467640.8286467820.8286467650.000000014

0.45	0.792606371	0.792606388	0.792606372	0.000000010	0.000000150
0.50	0.755635547	0.755635562	0.755635547	0.0000000005	0.0000000150
0.55	0.718443074	0.718443087	0.718443074	0.000000008	0.0000000130
0.60	0.681750643	0.681750657	0.681750644	0.000000013	0.000000130
0.65	0.646296385	0.646296395	0.646296386	0.0000000005	0.000000094
0.70	0.612838519	0.612838521	0.612838518	0.000000014	0.000000034
0.75	0.582159131	0.582159137	0.582159133	0.000000014	0.0000000041
0.80	0.555068148	0.555068152	0.555068148	0.0000000001	0.0000000040
0.85	0.532407454	0.532407455	0.532407454	0.0000000006	0.000000008
0.90	0.515055312	0.515055311	0.515055310	0.000000013	0.0000000004
0.95	0.503931027	0.503931026	0.503931027	0.0000000001	0.000000002
1.00	0.499999999	0.499999997	0.500000000	0.000000012	0.000000026

**Table 2** The results of HAM, HPM and Numerical methods for  $f'(\eta)$  when

S = 0.4, M = 2, A = 1

η	HAM	HPM	NUM	Error of HAM	Error of HPM
0.00	0.000000000	0.000000000	0.000000000	0.0000000000	0.0000000000
0.05	-0.138163081	-0.138163033	-0.138163081	0.0000000006	0.000000486
0.10	-0.261642988	-0.26164292	-0.261642989	0.000000014	0.000000688
0.15	-0.370750011	-0.370749941	-0.370750012	0.000000017	0.000000710
0.20	-0.465725958	-0.465725897	-0.465725960	0.000000023	0.000000626
0.25	-0.546745711	-0.546745665	-0.546745715	0.000000031	0.000000494
0.30	-0.613918173	-0.613918144	-0.613918177	0.0000000044	0.000000336
0.35	-0.667286685	-0.667286672	-0.667286691	0.000000061	0.000000184
0.40	-0.706828989	-0.706828991	-0.706828995	0.000000058	0.000000042
0.45	-0.732456763	-0.732456778	-0.732456770	0.000000072	0.000000085
0.50	-0.744014778	-0.744014806	-0.744014785	0.000000075	0.000000205
0.55	-0.741279668	-0.741279705	-0.741279676	0.000000086	0.000000288
0.60	-0.723958328	-0.723958365	-0.723958332	0.000000045	0.000000332
0.65	-0.691685870	-0.691685913	-0.691685876	0.000000054	0.000000373
0.70	-0.644023180	-0.64402322	-0.644023183	0.000000034	0.000000364
0.75	-0.580453874	-0.580453911	-0.580453878	0.000000035	0.000000330
0.80	-0.500380689	-0.500380723	-0.500380693	0.000000036	0.000000302
0.85	-0.403121088	-0.403121112	-0.403121091	0.000000033	0.000000210
0.90	-0.287901961	-0.287901979	-0.287901964	0.000000035	0.000000147
0.95	-0.153853214	-0.153853224	-0.153853218	0.0000000040	0.000000062
1.00	0.0000000030	0.0000000149	0.000000000	0.000000003	0.000000014

In this study, the problem of magneto-hydrodynamic MHD squeeze flow of an electrically conducting fluid between two infinite, parallel disks was analyzed using HAM and HPM. By the drawing of 2-D fig. 3 to 6, of the numerical solution, HPM and HAM solutions for f(y) and f'(y) with different values of S, M and A, we see that the Homotopy Analysis Method and Homotopy Perturbation Method are more accurate than NUM. According to fig. 3 to 6 and tables 1 and 2 these methods provide highly accurate analytic solutions for nonlinear problems in comparison with the numerical solution. The comparison of the method reveals that the approximations obtained by HAM and HPM converge to the exact solution quite fast. Also, the

auxiliary parameter  $\hbar$  provides us with a convenient way to adjust and control the convergence and its rate for the solutions series. Finally, it has been attempted to show the capabilities and wide-range applications of the HAM and HPM in comparison with the numerical solution of nonlinear equations.

### 7. CONCLUSION

In this paper, the problem of magneto-hydrodynamic MHD squeeze flow of an electrically conducting fluid between two infinite, parallel disks was analyzed by using HAM and HPM. Furthermore, the obtained solutions by HAM and HPM are compared with numerical solutions. The results demonstrate that HAM and HPM are very effective and simple.

They offer superior accuracy in comparison with the NUM. Also, it is found that these methods are powerful mathematical tools and that they can be applied to a large class of linear and nonlinear problems arising in different fields of science and engineering specially some heat transfer equations.

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